

Stern- und
Planetenentstehung
Sommersemester 2020
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Lecture 5: Stellar Feedback



http://exp-astro.physik.uni-frankfurt.de/star_formation/index.php

VORLESUNG/LECTURE

Raum: Physik - 02.201a

dienstags, 12:00 - 14:00 Uhr

SPRECHSTUNDE:

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dienstags: 14:00-16:00 Uhr

Nr.	Thema	Termin
1	Observing the cold ISM	21.04.2020
2	Observing Young Stars	28.04.2020
3	Gas Flows and Turbulence Magnetic Fields and Magnetized Turbulence	05.05.2020
4	Gravitational Instability and Collapse	12.05.2020
5	Stellar Feedback	19.05.2020
6	Giant Molecular Clouds	26.05.2020
7	Star Formation Rate at Galactic Scales	02.06.2020
8	Stellar Clustering	09.06.2020
9	Initial Mass Function – Observations and Theory	16.06.2020
10	Massive Star Formation	23.06.2020
11	Protostellar disks and outflows – observations and theory	30.06.2020
12	Protostar Formation and Evolution	07.07.2020
13	Late Stage stars and disks – planet formation	14.07.2020

5 STELLAR FEEDBACK

5.1 GENERAL FORMALISM

5.1.1 IMF –Averaged Yields

Instead of single stars we will describe populations of stars, i.e. collective properties of those populations.

IMF:
$$\xi(m) = \frac{dn}{d \ln m} \quad (\int \xi(m) dm = 1)$$

mean stellar mass:
$$\bar{m} = \frac{\int_{-\infty}^{\infty} m \xi(m) d \ln m}{\int_{-\infty}^{\infty} \xi(m) d \ln m} = \frac{1}{\int_{-\infty}^{\infty} \xi(m) d \ln m} = \frac{\text{total mass}}{\text{total number}}$$

Q : quantity produced by stars with rate q (e.g: Q : total radiant energy E , q : bolometric luminosity $L(m, t)$ of a star of mass m and age t)

Instantaneous production rate:

$$q(t) = M \int_{-\infty}^{\infty} d \ln m \xi(m) q(m, t)$$

IMF-averaged production rate:

$$\left\langle \frac{q}{M} \right\rangle = \int_{-\infty}^{\infty} d \ln m \xi(m) q(m, t)$$

$\langle q/M \rangle$: instantaneous rate at which stars are producing Q per unit stellar mass

Over lifetime of stellar population:

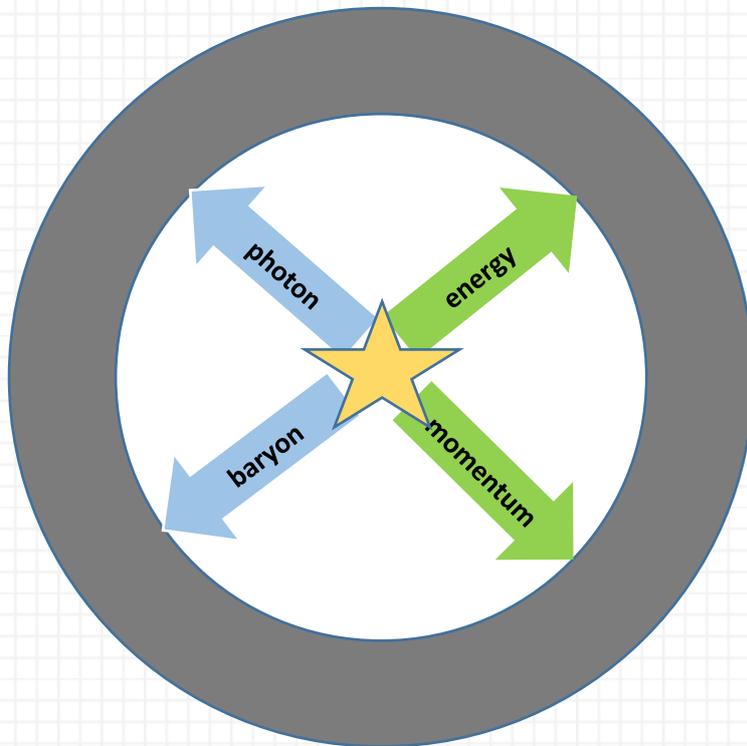
$$Q = M \int_{-\infty}^{\infty} d \ln m \xi(m) \int_0^{\infty} dt q(M, t)$$

IMF-averaged yield:

$$\left\langle \frac{Q}{M} \right\rangle = \int_{-\infty}^{\infty} d \ln m \xi(m) \int_0^{\infty} dt q(M, t)$$

$\langle Q/M \rangle$: total amount of Q produced per unit stellar mass of stars formed over the stars' entire lifetimes.

5.1.2 Energy- versus Momentum-Driven Feedback



stars inject energy and momentum into the surrounding ISM (by photons and baryons)

gas heating difficult because energy can be radiated away

energy input (e.g. by FUV radiation) only important if mechanism can keep the ISM hot for a significant time (e.g. crossing time)

momentum-driven feedback: feedback mechanism where energy is lost rapidly

energy-driven feedback: feedback mechanism where energy is kept for some significant time

Consider a star surrounded by a uniform region of gas, that turns on at $t = 0$ with energy injection rate \dot{E}_w and momentum injection rate \dot{p}_w (by photons and baryons, i.e. stellar wind)

Properties of the gas shell:

all energy radiated away

momentum of shell at t is all the radial momentum injected up to t .

$$p_{sh} = M_{sh} v_{sh} = \dot{p}_w t$$

kinetic energy of the shell ($v_{sh} \sim \text{few } 10 \text{ km s}^{-1}$)

$$E = \frac{p_{sh}^2}{2M_{sh}} = \frac{1}{2} v_{sh} \dot{p}_w t$$

no energy radiated away

kinetic energy

$$E = \dot{E}_w t$$

$$\frac{\text{E in energy - conserving case}}{\text{E in all energy - lost case}} = \frac{1}{v_{sh}} \frac{2\dot{E}_w}{\dot{p}_w}$$

If \dot{E}_w is carried by baryons $\left[\frac{2\dot{E}_w}{\dot{p}_w} = v_{wind} \right]$ ($v_{wind} \sim \text{few } 1000 \text{ km s}^{-1}$)

If \dot{E}_w is carried by photons $\left[\frac{2\dot{E}_w}{\dot{p}_w} = 2c \right]$

It matters a great deal whether the energy is kept in the ISM or not!

5.2 MOMENTUM-DRIVEN FEEDBACK MECHANISMS

5.2.1 Radiation Pressure & Radiatively-Driven Winds

radiation pressure is probably momentum-driven (because radiation energy will be re-radiated quickly in most cases)

Murray & Rahman (2010) find

$$\left\langle \frac{L}{M} \right\rangle = 1140 L_{\odot} M_{\odot}^{-1} = 2200 \text{ erg g}^{-1}$$

Then the momentum injection rate is

$$\left\langle \frac{p_{rad}}{M} \right\rangle = \frac{1}{c} \left\langle \frac{L}{M} \right\rangle = 7.3 \times 10^{-8} \text{ cm s}^{-2} = 23 \text{ km s}^{-1} \text{ Myr}^{-1}$$

Every gram of matter that goes into stars, those stars produce enough light over 1 Myr to accelerate another gram of ISM matter to a speed of 23 km s⁻¹.

If massive stars produce strong winds, they might carry ~ 50% of the momentum of the photons => $\left\langle \frac{p_{rad}}{M} \right\rangle$ becomes larger.

Total energy production (integrated over 100 Myr)

$$\left\langle \frac{E_{rad}}{M} \right\rangle = 1.1 \times 10^{51} \text{ erg } M_{\odot}^{-1}$$

$$L \sim M^{3.5}$$

Majority produced in first 5 Myr (by massive stars).

Energy budget in units of c^2 gives dimensionless efficiency by which stars convert mass into energy:

$$\epsilon = \frac{1}{c^2} \left\langle \frac{E_{rad}}{M} \right\rangle = 6.2 \times 10^{-4}$$

Radiation momentum budget:

$$\left\langle \frac{p_{rad,tot}}{M} \right\rangle = \frac{\epsilon}{c} = 190 \text{ km s}^{-1}$$

Total radiation momentum output by stars is significant (compare with circular velocity of a typical galaxy ~200 km s⁻¹).

5.2.2 Protostellar Winds



Abbildung 1 Credit: NASA/JPL-Caltech/Univ. of Michigan, Spitzer image: Infrared light with a wavelength of 3.6 microns has been color-coded blue; 4.5-micron light is green; and 8.0-micron light is red.

Accretion disks produce a wind (jet) that carries away some of the accreted mass and angular momentum (few 10% of the total accreted mass) at a velocity \sim Keplerian speed at the stellar surface.

Consider a star of mass M_* and radius R_* . Its wind will move at a speed of order

$$v_w \sim \sqrt{G \frac{M_*}{R_*}} = 250 \text{ km s}^{-1} \left(\frac{M_*}{M_\odot} \right)^{\frac{1}{2}} \left(\frac{R_*}{3R_\odot} \right)^{-1/2}$$

The kinetic energy per unit mass carried by the wind is $v_w^2/2$ and when the wind hits the surrounding ISM it will shock and this kinetic energy will be converted to thermal energy. We can therefore find the post-shock temperature from energy conservation. The thermal energy per unit mass is $\left(\frac{3}{2}\right) k_B T / \mu m_H$ thus the post-shock temperature will be

$$T = \frac{\mu m_H v_w^2}{3k_B} \sim 5 \times 10^6 K$$

This is low enough that gas at this temperature will be able to cool fairly rapidly, therefore proto-stellar winds are most-likely momentum-driven because (shocked gas ($T \sim$ few million K) cools quickly via radiation)

Consider stars forming over a time t_{form} with constant accretion rate.

Thus, a star of mass m accretes with a rate $\dot{m} = \frac{m}{t_{form}}$, over the forming time and during that time it produces a wind with a mass flux $f\dot{m}$ that is launched at speed v_K .

IMF-averaged yield of wind momentum:

$$\left\langle \frac{p_w}{M} \right\rangle = \int_{-\infty}^{\infty} d \ln m \xi(m) \int_0^{t_{form}} dt \frac{f m v_K}{t_{form}} = f v_K \int_{-\infty}^{\infty} d \ln m \xi(m) = f v_K$$

$$f v_K \sim \text{few } 10 \text{ km s}^{-1}$$

assuming f and v_K are constant

- ⇒ protostellar winds supply over their full lifetime as much momentum as radiation does in 1 Myr.
- ⇒ Compared to total stellar lifetime radiation injects much more momentum!
- ⇒ But radiation pressure is strongest from most massive stars, but most stars are low mass stars.
- ⇒ Protostellar winds produce the \sim same amount of momentum independent of stellar mass, every star contributes the same
- ⇒ During the accretion time, winds are much more powerful than radiation

5.3 (PARTLY) ENERGY-DRIVEN FEEDBACK MECHANISMS

5.3.1 Ionizing Radiation

$$h\nu > 13.6 \text{ eV}$$

Massive stars produce ionizing radiation. Murray & Rahman (2010) find a yield of ionizing photons from a zero-age population:

$$\left\langle \frac{S}{M} \right\rangle = 6.3 \times 10^{46} \text{ photons s}^{-1} M_{\odot}^{-1}$$

Averaged over total lifetime

$$\left\langle \frac{S_{tot}}{M} \right\rangle = 4.2 \times 10^{60} \text{ photons } M_{\odot}^{-1}$$

5.3.1.1 HII Region Expansion

$\text{H II} \leftarrow$ singly ionized
 $\text{I} \leftarrow$ neutral

(Hydrogen) Ionizing photons will be absorbed with a very short mean free path, producing a bubble of H^+ gas within which all photons are absorbed.

The size of the bubble (Strömgen radius) results from equating the recombination rate with the production rate of ionizing photons

$$S = \frac{4}{3} \pi r_S^3 n_e n_p \alpha_B$$

$$\alpha_B \approx 3 \times 10^{-13} \text{ cm}^3 \text{ s}^{-1}$$

α_B : recombination rate

r_S : Strömgen radius

n_e, n_p : electron and proton number densities

μ_H : mean mass per hydrogen nucleus in the gas

ρ_0 : initial density before photoionization starts

$n_p = \rho_0 / \mu_H$ and $n_e = 1.1 \rho_0 / \mu_H$ (1.1 because we add singly ionized He 1/10 He per H nucleus)

$$r_S = \left(\frac{3S\mu_H^2}{4(1.1)\pi\alpha_B\rho_0^2} \right)^{1/3} = 2.8 S_{49}^{1/3} n_2^{-2/3} \text{ pc}$$

$$S_{49} = S/10^{49} \text{ s}^{-1}, n_2 = \frac{(\rho_0/\mu_H)}{100 \text{ cm}^{-3}}, \alpha_B = 3.46 \times 10^{-13} \text{ cm}^3 \text{ s}^{-1}$$

The ionized gas will be heated to $\sim 10^4$ K. The corresponding sound speed is

$$c_i = \sqrt{\frac{k_B T_i}{\mu_H/2.2}} = 11 T_{i,4}^{1/2} \text{ km s}^{-1}$$

where $T_{i,4} = T_i/10^4\text{K}$, $\mu_H = 2.3 \times 10^{-24} \text{ g}$ (2.2 because that is the number of particles per H nucleus, i.e. electrons)

Pressure inside r_s $P_i = \rho_0 c_i^2 \gg \rho_0 c_0^2 = P_{ISM}$ ambient pressure

⇒ HII region will start to expand

⇒ time scale for ionization balance is short compared to dynamical time scale ⇒ ionization balance will remain during expansion

For a given radius r_i after expansion started, the density inside the HII region will be

$$\rho_i = \left(\frac{3S\mu_H^2}{4(1.1)\pi\alpha_B r_i^3} \right)^{1/2}$$

⇒ The expanding HII region sweeps up the surrounding gas

⇒ mass accumulates in a neutral shell surrounding the HII region

⇒ at late times, when $r_i \gg r_s$, and $\rho_i \ll \rho_0$ we can neglect the mass in the shell interior

⇒ then the shell mass is $M_{sh} = \frac{4}{3}\pi\rho_0 r_i^3$

Equation of motion for shell (neglecting ambient pressure):

$$\frac{d}{dt}(M_{sh}\dot{r}_i) = 4\pi r_i^2 \rho_i c_i^2$$

$$\frac{d}{dt}\left(\frac{1}{3}r_i^3 \dot{r}_i\right) = c_i^2 r_i^2 \left(\frac{r_i}{r_s}\right)^{-3/2}$$

$$\rho_i = \rho_0 \left(\frac{r_i}{r_s}\right)^{-3/2}$$

General solution: numerically, for late times, when $r_i \gg r_s$ analytic solution possible

If $r_i \gg r_s$ we can take $r_i(t \rightarrow 0) \approx 0$ as BC

→ similarity solution of the ODE

trial solution $r_i = f r_s \left(\frac{t}{t_s}\right)^\eta$, with $t_s = \frac{r_s}{c_i} = 240 S_{49}^{\frac{1}{3}} n_2^{-\frac{2}{3}} T_{i,4}^{-\frac{1}{2}} \text{ kyr}$

insert r_i

$$\frac{1}{4} \eta (4\eta - 1) f^4 \left(\frac{t}{t_s}\right)^{4\eta-2} = f^{1/2} \left(\frac{t}{t_s}\right)^{\eta/2}$$

solution only if $4\eta - 2 = \eta/2$, therefore $\eta = 4/7$, and $f = \left(\frac{49}{12}\right)^{2/7}$

$$r_i = r_s \left(\frac{7t}{2\sqrt{3}t_s}\right)^{\frac{4}{7}} = 9.4 S_{49}^{\frac{1}{7}} n_2^{-\frac{2}{7}} T_{i,4}^{\frac{2}{7}} t_6^{\frac{4}{7}} \text{ pc}$$

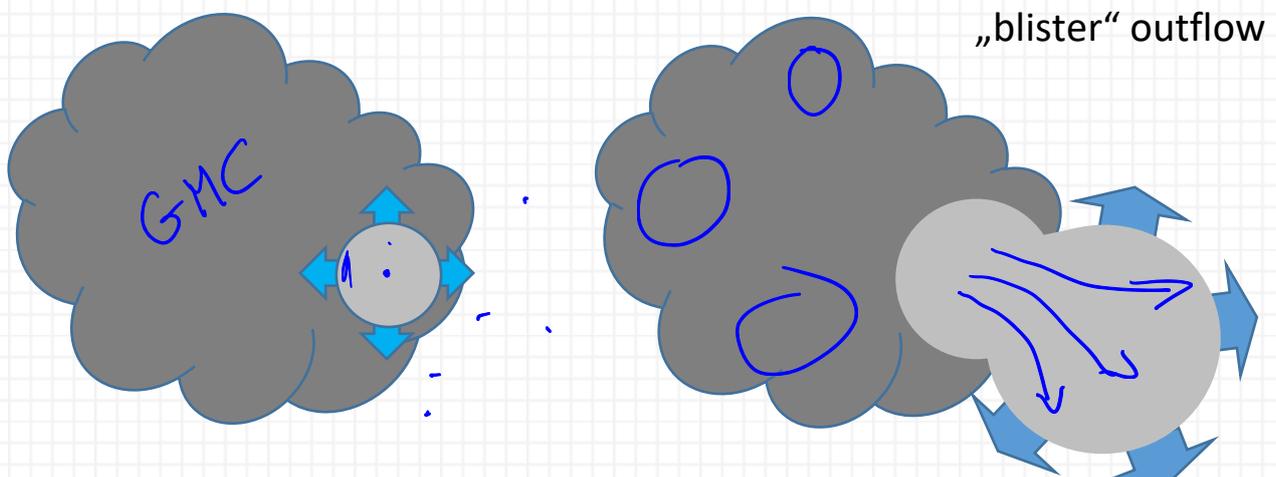
at late times, where $t_6 = t/1\text{Myr}$

5.3.1.2 Feedback Effects of HII Regions

Ionization can

- ⇒ eject mass
- ⇒ drive turbulent motion
- ⇒ possibly disrupt clouds

Mass ejection:



Maximum mass flux carried by the ionized wind:

$$\dot{M} = 4\pi r_i^2 \rho_i c_i$$

$$\dot{M} = 4\pi r_S^2 \rho_0 c_i \left(\frac{7t}{2\sqrt{3}t_S} \right)^{\frac{2}{7}} = 7.2 \times 10^{-3} t_6^{2/7} S_{49}^{4/7} n_2^{-1/7} T_{i,4}^{1/7} M_\odot \text{yr}^{-1}$$

Over the lifetime of an O star: 3-4 Myr, it can eject $\sim 10^3 - 10^4 M_\odot$ of mass from its parent cloud (if n_2 is not too high)

Massive stars can eject many times their own mass from a GMC.

Energy contained in the expanding shell

$$E_{sh} = \frac{1}{2} M_{sh} \dot{r}_i^2 = \frac{32}{147} \pi \rho_0 \frac{r_S^5}{t_S^2} \left(\frac{7t}{2\sqrt{3}t_S} \right)^{6/7}$$

$$= 8.1 \times 10^{47} t_6^{6/7} S_{49}^{5/7} n_2^{-10/7} \text{erg}$$

Compare: binding energy of $10^5 M_\odot$ GMC with surface density of 0.03 g cm^{-2} is $\sim 10^{50} \text{ erg}$.

⇒ single O star has a comparable little effect on the GMC

⇒ cluster of ~ 100 O stars can produce HII regions whose energies are comparable to the total binding energy of the GMC

⇒ sometimes HII regions can disrupt entire GMCs

Momentum of the expanding shell:

$$p_{sh} = M_{sh} \dot{r}_i = 1.1 \times 10^5 n_2^{-1/7} T_{i,4}^{-8/7} S_{49}^{4/7} t_6^{9/7} M_\odot \text{km s}^{-1}$$

Non-linearity in S_{49} and time \Rightarrow effects depend on how the stars are clustered and how long they live

Typical numbers: $L_{ionization} = 10^{49}$ photons/s, age 4 Myr,
 $n_2 = 1$, and $T_{i,4} = 1$

momentum injected per 10^{49} photons s^{-1}

$$\Rightarrow p = 3 - 5 \times 10^5 M_{\odot} \text{ km s}^{-1}$$

A zero-age population gives 6.3×10^{46} photons $s^{-1} M_{\odot}^{-1}$

The momentum injection rate for HII regions is (very roughly)

$$\left\langle \frac{\dot{p}_{HII}}{M} \right\rangle \sim 3 \times 10^3 \text{ km s}^{-1}$$

HII regions are most likely the dominant feedback.

Limitation: $v_i(t=0) \sim 10 \text{ km s}^{-1}$, very massive cluster might have comparable escape velocities, it will prevent disruption of the GMC

5.3.2 Stellar Winds

O stars have winds of $v_w \sim 1000 - 2500 \text{ km s}^{-1}$ and mass fluxes of $\dot{M}_w \sim 10^7 M_{\odot} \text{ yr}^{-1}$.

⇒ Momentum small compared to other effects.

⇒ high v_w produces very high post-shock temperatures of $T \sim 10^8 \text{ K}$

⇒ gas with $T \sim 10^8 \text{ K}$ has a very long cooling time

⇒ **energy-driven feedback**

Wind momentum comparable to stellar radiation momentum (Repolust et al. 2004):

$$\dot{M}_w v_w \approx 0.5 \frac{L_*}{c}$$

L_* : stellar luminosity

Mechanical luminosity of the wind

$$L_w = \frac{1}{2} \dot{M}_w v_w^2 = \frac{L_*^2}{8 \dot{M}_w c^2} = 850 L_{*,5}^2 \dot{M}_{w,-7}^{-1} L_\odot$$

The integrated power output over the ~ 4 Myr lifetime of a massive star

$$E_w = L_w t = 1.0 \times 10^{50} L_{*,5}^2 \dot{M}_{w,-7}^{-1} t_6 \text{ erg}$$

The total mechanical power in the wind is comparable to the amount of energy released when the star goes supernova.

Equation of motion: assume energy is conserved (1/2 goes into heating)

$$\frac{d}{dt} \left(\frac{2}{3} \pi \rho_0 r_b^3 \dot{r}_b^2 \right) \approx \frac{1}{2} L_w$$

Similar to HII region we try a similarity solution with:

$$r_b = A t^\eta$$

$$\frac{4}{3} \pi \eta^2 (5\eta - 2) \rho_0 A^5 t^{5\eta-3} \approx L_w$$

Because of energy conservation this must be independent of t and then implies $\eta = 3/5$ and $A = [25 L_w / (12 \pi \rho_0)]^{1/5}$.

$$r_b = 16 L_{*,5}^{2/5} \dot{M}_{w,-7}^{-1/5} n_2^{-1/5} t_6^{3/5} \text{ pc}$$

(greater than radius of HII region!)

- ⇒ Wind will move faster and drive the ionized gas into a thin layer between hot wind gas and outer cool shell – *if the energy-driven limit is correct!*
- ⇒ Major assumption that energy remains confined in a closed shell may not be true. Possibly break out of gas.
- ⇒ Theoretically difficult to predict.
- ⇒ Observations can help: if shocked gas is trapped it should produce x-ray emission!

Pressure of x-ray emitting gas:

Shell expanding for time t . Then the total energy within that shell is $E_w \approx L_w t$.

Pressure is 2/3 of energy density (monoatomic gas)

$$P_X = \frac{2E_W}{3[(4/3)\pi r^3]} = \frac{L_* t}{16\pi \dot{M}_w c^2 r^3}$$

The pressure exerted by the radiation (twice that exerted by the wind in the momentum-driven limit)

$$P_{rad} = \frac{L_*}{4\pi r^2 c}$$

ratio is called trapping factor:

$$f_{trap} = \frac{P_X}{P_{rad}} = \frac{L_* t}{4\dot{M}_w c r} \approx \frac{L_*}{4\dot{M}_w c v}$$

$$v = r/t$$

expansion vel. of shell

Using $\dot{M}_w v_w \approx \left(\frac{1}{2}\right) L_* / c$

$$f_{trap} = \frac{v_w}{2v}$$

If the shell expands in energy-driven limit due to winds: the pressure of the hot gas within should exceed the direct radiation pressure by about a factor v_w/v .

In the momentum-driven limit $P_X/P_{rad} \sim 1/2$

Example: Lopez et al. 2010 observed 30 Doradus

Theory: $v_{expansion} \approx 20 \text{ km s}^{-1}$, $v_w \approx 1000 \text{ km s}^{-1} \Rightarrow f_{trap} = 20$

Obs: $f_{trap} \approx 0.5$ therefore momentum-driven limit more likely.

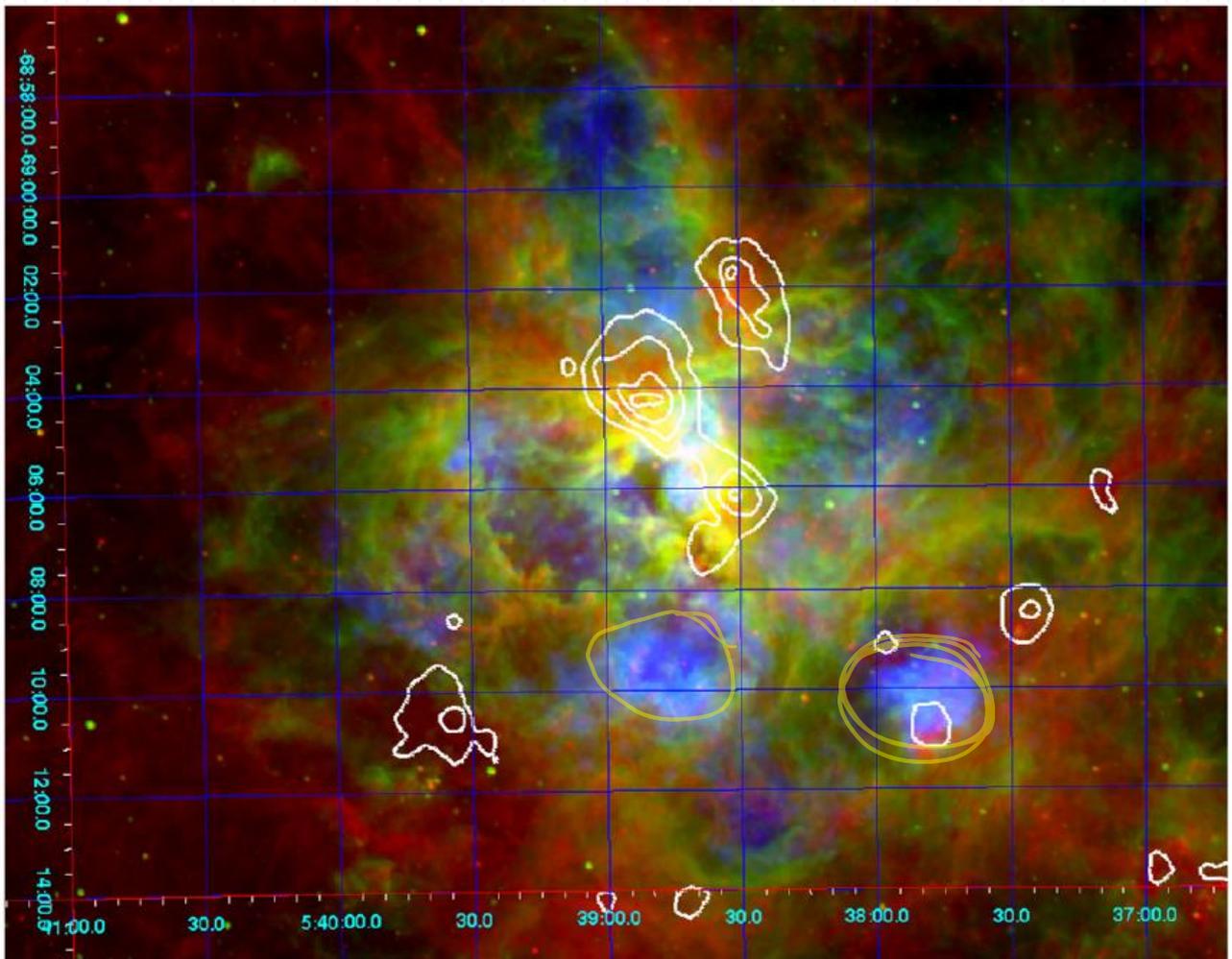


Abbildung 2 Lopez et al. 2012 : Three-color image of 30 Doradus: MIPS 8 μm (red), H α (green), and 0.5–8 keV X-rays (blue). White contours show the 12CO(1–0) emission (Johansson et al. 1998) in the region. Both large- and small-scale structures are evident. north is up, east is left.

5.3.3 Supernovae

Model $q(m, t)$ as δ function: all the energy and momentum of the SN released in a single burst at a time $t = t_l(m)$, where $t_l(m)$ is the lifetime of the progenitor star.

We assume an energy yield per star of 10^{51} erg. Its minimum mass is $\sim 8M_{\odot}$

SN energy per unit mass:

$$\left\langle \frac{E_{SN}}{M} \right\rangle = E_{SN} \int_{m_{min}}^{\infty} d \ln m \xi(m) = E_{SN} \left\langle \frac{N_{SN}}{M} \right\rangle$$

$\langle N_{SN}/M \rangle$: number of stars above $8M_{\odot}$, which is the number of expected SN

For the Chabrier IMF from $0.01 - 120M_{\odot}$ we find:

$$\left\langle \frac{N_{SN}}{M} \right\rangle = 0.011 M_{\odot}^{-1}, \quad \left\langle \frac{E_{SN}}{M} \right\rangle = 1.1 \times 10^{49} \text{erg} M_{\odot}^{-1} = 6.1 \times 10^{-6} c^2$$

$$\text{MW SFR} \sim 1 M_{\odot} \text{yr}^{-1} \Rightarrow \underline{1 \text{ SN per century in the MW}}$$

Momentum yield from SN:

ejection velocity: $v_e \approx 10^9 \text{cm s}^{-1}$

momentum: $p_{SN} = \underline{2E_{SN}/v_{ej}}$

$$\left\langle \frac{p_{SN}}{M} \right\rangle = \frac{2}{v_{ej}} \left\langle \frac{E_{SN}}{M} \right\rangle = 55 \underline{v_{ej,9}^{-1}} \text{ km s}^{-1}$$

(small compared to other effects)

