

# A new storage ring for ISOLDE

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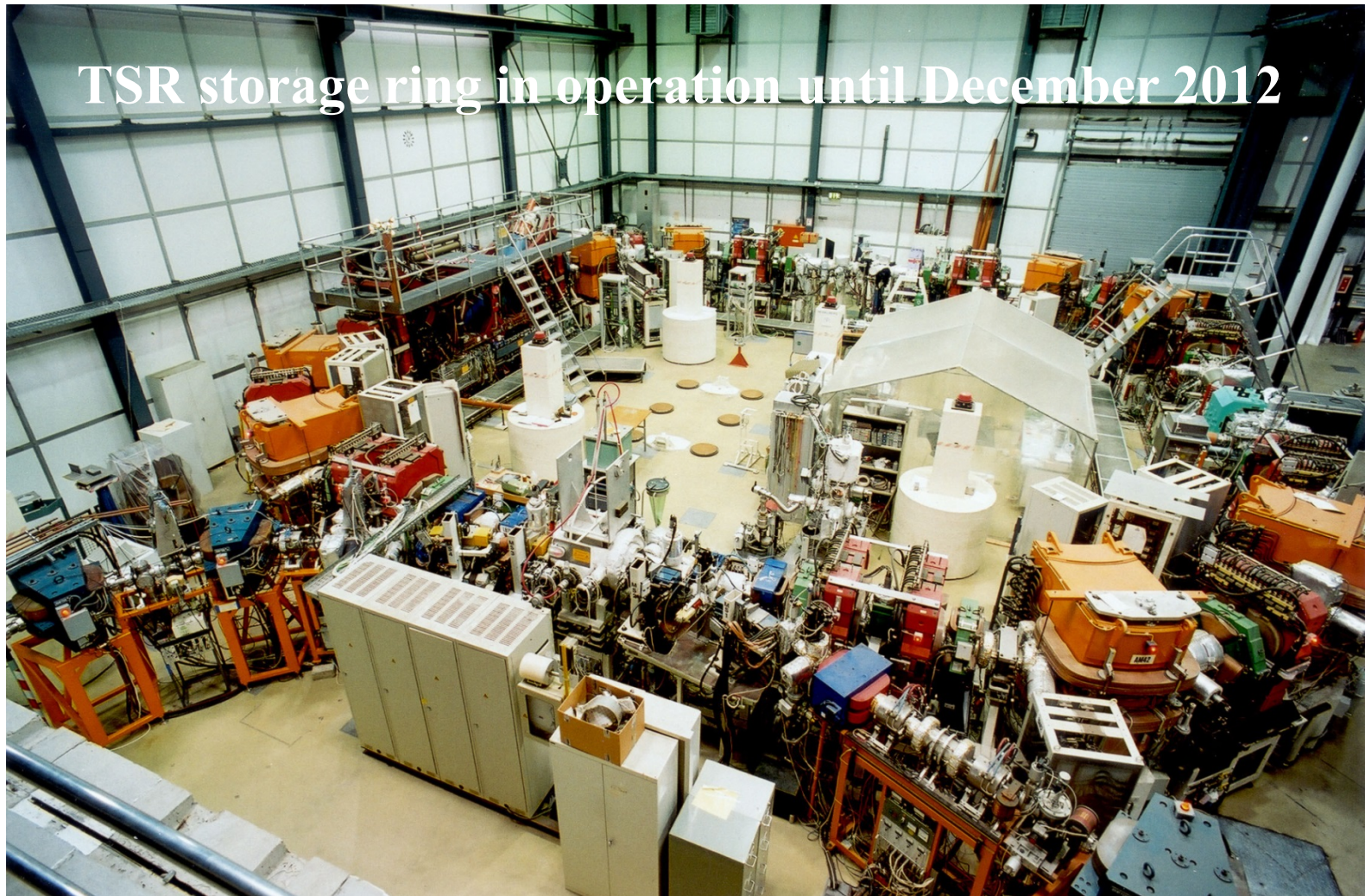
Max Planck Institut für Kernphysik, Heidelberg



NARRS workshop, GSI, Darmstadt, 13<sup>th</sup>-15<sup>th</sup> March 2018

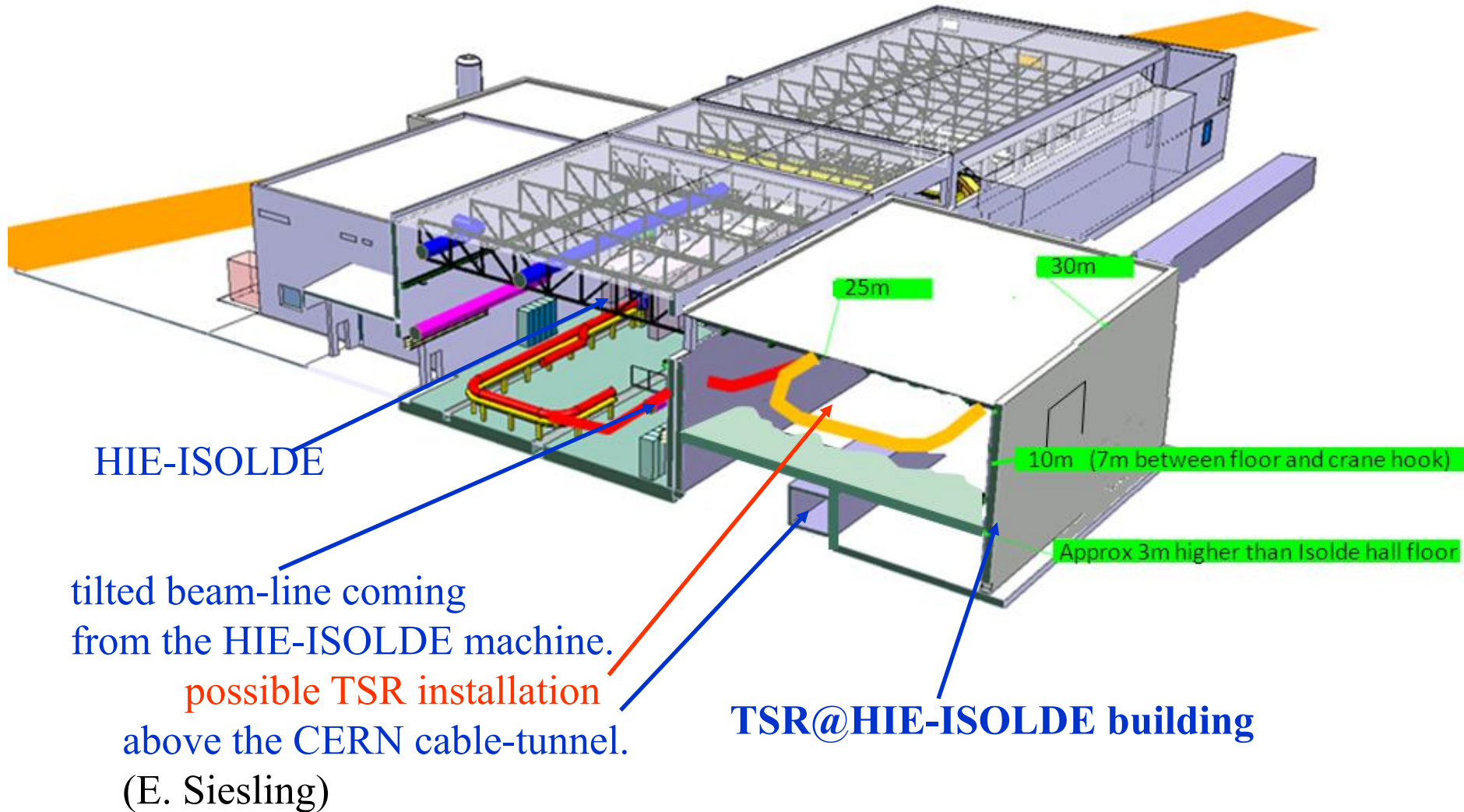
# Proposed TSR@ISOLDE project

to store radioactive ions for nuclear physics experiments it was proposed to move TSR located at MPI for nuclear physics to ISOLDE.





# TSR @ HIE-ISOLDE



# Time-line of the TSR@ISOLDE project

TSR@ISOLDE workshop at MPI-K Heidelberg  
evaluated the future for TSR **Oct 2010**

ISOLDE and Neutron Time-of-Flight Committee endorsed **Jan 2012**

TSR technical design report **129 co-authors (47 institutions)**

EPJ Special Topics **207 1-117 May 2012** →

Approved by CERN Research board, **May 2012**

*“The installation of TSR, as an experiment to be included in the HIE-ISOLDE programme, was approved by the Research Board.*

*The timescale will be defined once the study of its Integration has been completed.”*

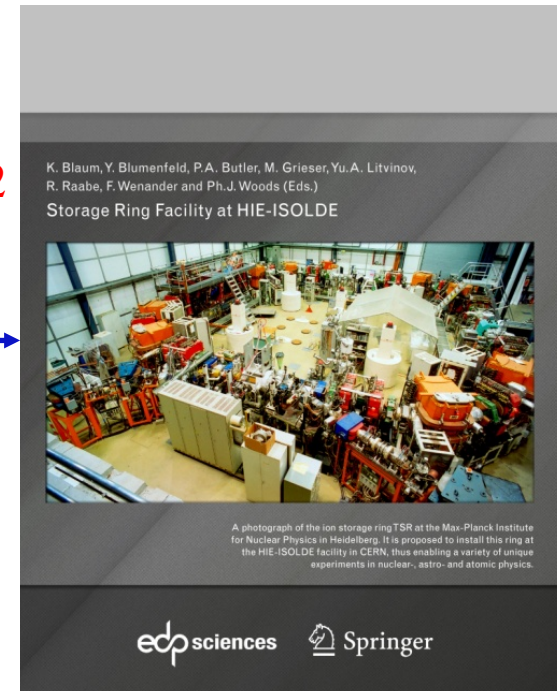
Presentation of the integration study to the CERN Research Board **Nov 2013**

Several TSR@ISOLDE workshops at CERN: **2012, 2014, 2015**

Updated CERN integration study with report to the CERN directorate **2016**

CERN director general: decision about the TSR@ISOLDE project is postponed until 2020/2021 (after second LHC upgrade ) **August/September 2016**

MPIK cannot hold TSR at MPIK until 2020/2021 without getting green light from CERN



# Design Criteria of the new storage ring

a) storage ring should be able to store ions up to  $^{238}\text{U}$  and 10 MeV/u at the equilibrium charge state obtainable with the HIE-ISOLDE stripper.

⇒ maximum rigidity of the ring:  $B\rho_{\text{max}} \approx 1.5 \text{ Tm}$

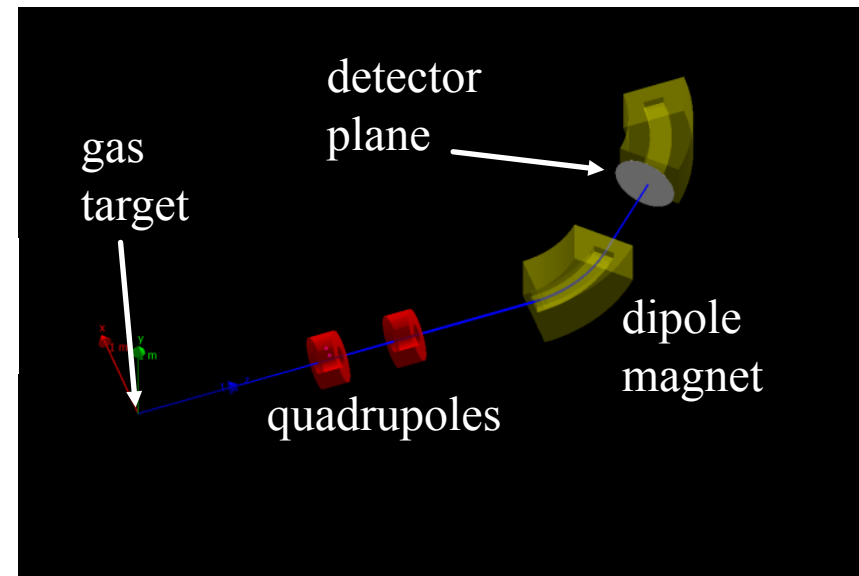
b) storing of heavy daughter nuclei up to a certain rigidity deviation ( $\Delta B\rho/B\rho$ ) created in nuclear reactions should be possible.

⇒ small maximum dispersion function  $D_x$  of the storage ring

dispersion at the gas jet  $D_x=0 \text{ m}$  to avoid a excitation of betatron oscillation of daughter nuclei

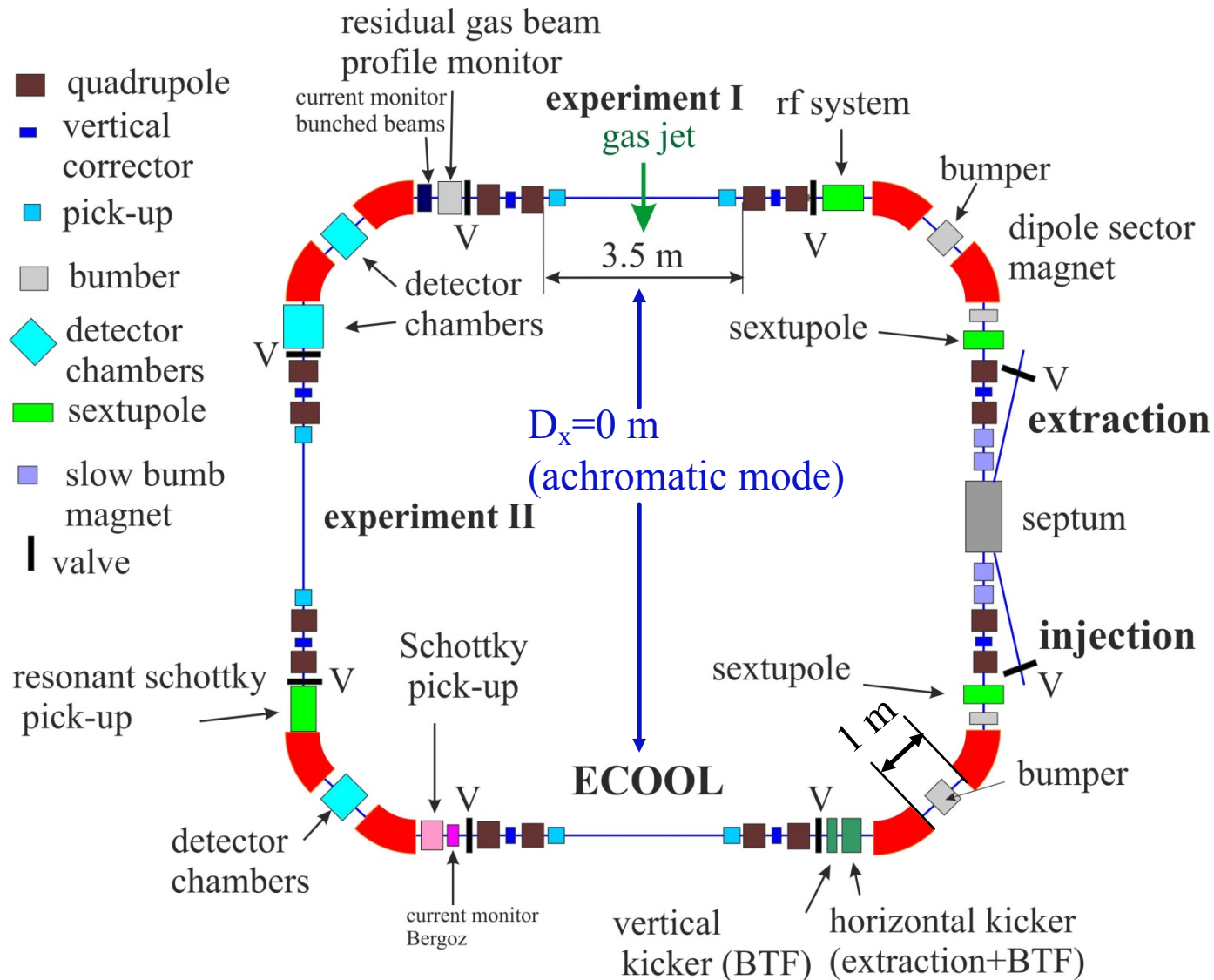
c.) daughter nuclides with large transfers momenta produced in nuclear reactions should be focused at the detector position

d.) storing ring should be compact to fit in the present HIE Isolde hall



# Layout of the new storage ring

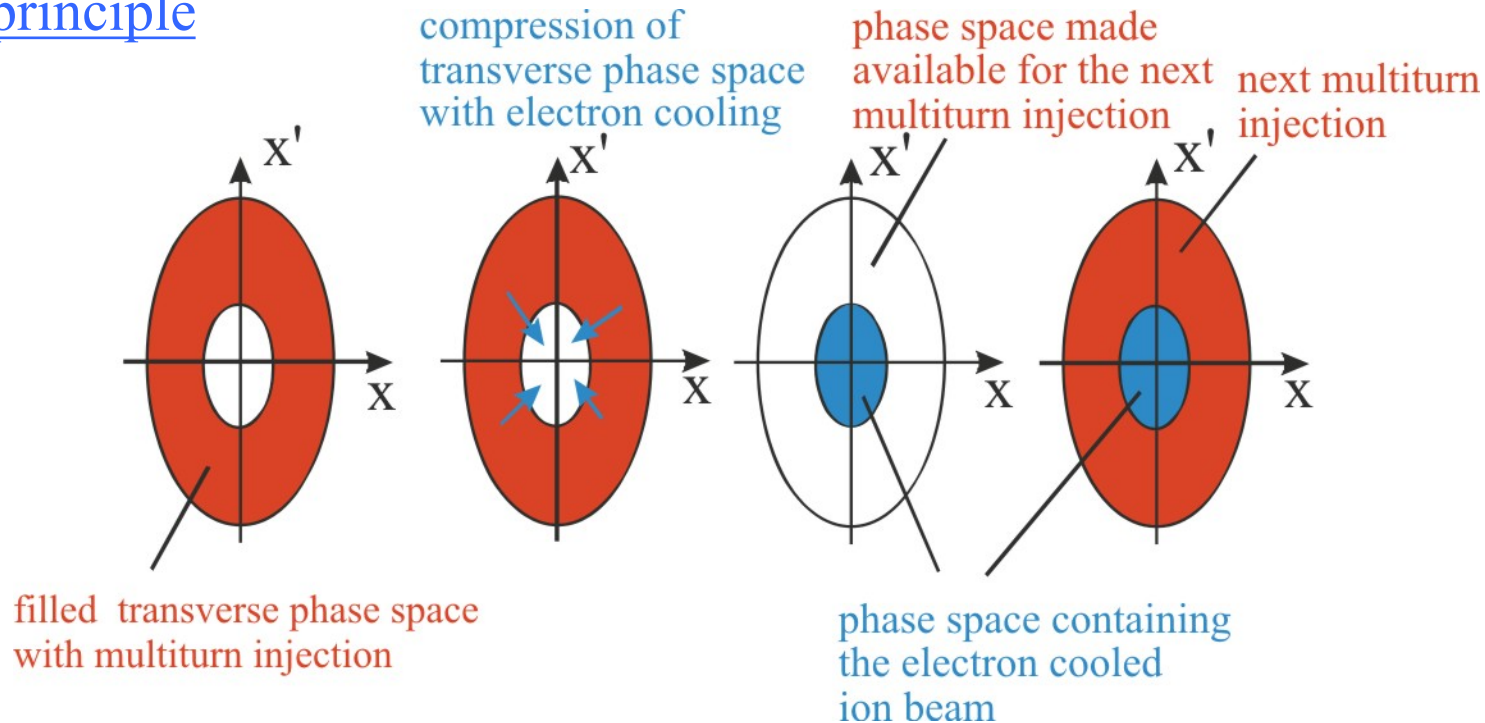
with straight section Length  $L=3.5$  m and circumference  $C=42.4$  m



# ECOOOL Stacking

A combination of multi-turn injection and electron cooling stacking will be used to fill the storage ring with particles

## principle



particle number  $N(t)$ :

$$\frac{dN(t)}{dt} = n_r N_i - \frac{N(t)}{\tau}$$

$N_i$ -injected particle number per injection

$n_r$ - injection rate

$\tau$ - total lifetime

# Luminosity with ECOOL stacking

In the equilibrium:  $dN(t)/dt = 0$   
 number of stored ions  $N_0$  ( $N_0 \leq N_s$ ):

$$N_0 = n_r \tau N_i$$

$N_0$  determines the luminosity  $L$ :

$$L = \frac{R}{\sigma} = N_0 f_0 n_t = n_r \tau n_t N_i f_0$$

$\uparrow$   
 $\tau \sim \frac{1}{n_t}$

$N_0$  - equilibrium particle number

$N_s$  - space charge limit

$R$  - reaction rate

$\sigma$  - cross section

$N_i$  - injected number of ions with  
 multi-turn injection

$n_t$  - target thickness

$n$  - target density  $n_t = \int n ds$

$f_0$  - revolution frequency

$n_r$  - injection rate  $n_r = 1/T$

$T$  - time between two injections

for large  $n_t$  the total life time  $\tau$  given by the target thickness  $n_t$ :

$$L = \frac{R}{\sigma} = \text{const} \frac{N_i}{T}$$

$\uparrow$   
 independent of  
 target thickness !

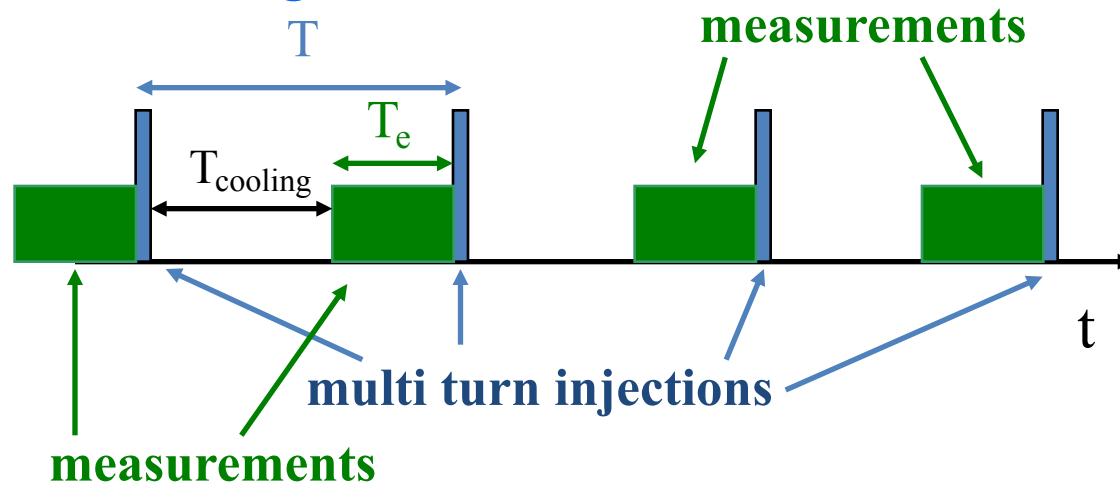
$\longleftarrow$  number of ions per multi turn injection  
 $\longleftarrow$  time between two multi turn injections



# Effective Luminosity

- continuous injection cycle
- measurement starts after the cooling time  $T_{\text{cooling}}$  of the injected ions and ends before the next multi-turn injection takes place

## injection and measuring scheme



Luminosity is reduced by factor  $\eta = T_e/T$   
 $\Rightarrow$  effective Luminosity  $L_{\text{eff}}$ :

$$L_{\text{eff}} = \frac{R}{\sigma} = \eta N_0 f_0 n_t = \eta \frac{\tau}{T} f_0 n_t N_i$$

↑

below space charge limit  $N_0 \leq N_s$      $N_s$ -space charge limit

$N_0$ -total number of stored particles  
 $N_i$ -injected particle number per multi-turn injection  
 $T$ - time between two injections  
 $\tau$ - total lifetime  
 $f_0$ -revolution frequency

# Space charge limit due to incoherent tune

maximum possible stored ion number: 
$$N_s = \frac{A}{q^2} \frac{2\pi}{r_p} \cdot B \cdot \beta^2 \cdot \gamma^3 \cdot \varepsilon \cdot (-\Delta Q)$$

$-\Delta Q$  - possible incoherent tune shift for  $B=1$  at TSR:  $-\Delta Q \approx 0.065 - 0.1$

for an electron cooled ion beam:

$$\varepsilon \propto \left( \frac{q^4}{A^2} \frac{N_s}{\lambda_{\text{Cool}}} \frac{1}{\beta^3} \right)^{0.44} \quad \lambda_{\text{cool}} \propto n_e \frac{q^2}{A} \quad n_e \propto \beta^2$$

new storage ring (ISR) has similar possible incoherent tune shifts

incoherent tune shift for an electron cooled ion beam

$$N_s = \text{const} \frac{(A^{33}/E^5)^{1/28}}{q^2}$$

E-ion energy in MeV

A-ion mass

q-ion charge state

TSR experiments:  $\text{const} \approx 7 \cdot 10^9$

**Measured space charge limit of an electron cooled ion beam at the TSR**

Ion	E (MeV)	measured $N_s$	calculated $N_s$
p	21	$5.4 \cdot 10^9$	$4.1 \cdot 10^9$
$^{16}\text{O}^{8+}$	98	$9.4 \cdot 10^8$	$1.3 \cdot 10^9$
$^{12}\text{C}^{6+}$	73	$1.7 \cdot 10^9$	$1.7 \cdot 10^9$
$^{32}\text{S}^{16+}$	195	$9.5 \cdot 10^8$	$6.3 \cdot 10^8$
$^{35}\text{Cl}^{17+}$	293	$5.1 \cdot 10^8$	$5.8 \cdot 10^8$

$\Rightarrow$  maximum possible Luminosity for  $N_0 = N_s$

# Effective Luminosity of some selected ions

$L_{\text{eff}} = \eta N_0 f_0 n_t$  calculated for the space charge limit ( $N_0 = N_s$ ) and a beam life time  $\tau = 1-1.5\text{s}$

beam	Energy (MeV/u)	$q_s$	$N_s$	target	$n_t$ (atoms/cm <sup>2</sup> )	$\tau$ (s)	$N_i/T$ (1/s)	$L_{\text{eff}}$ (1/cm <sup>2</sup> s)	$\eta = 0.5$
<sup>86</sup> Kr <sup>36+</sup>	10	36+	$3 \cdot 10^8$	H <sub>2</sub>	$9 \cdot 10^{14}$	1.5	$3 \cdot 10^8$	$2 \cdot 10^{29}$	proton capture reactions
<sup>96</sup> Ru <sup>39+</sup>	10	39+	$3 \cdot 10^8$	H <sub>2</sub>	$4 \cdot 10^{14}$	1	$3 \cdot 10^8$	$8 \cdot 10^{28}$	
<sup>196</sup> Hg <sup>64+</sup>	10	64+	$2 \cdot 10^8$	H <sub>2</sub>	$5 \cdot 10^{13}$	1.5	$2 \cdot 10^8$	$8 \cdot 10^{27}$	
<sup>232</sup> Th <sup>71+</sup>	10	71+	$2 \cdot 10^8$	H <sub>2</sub> , D <sub>2</sub>	$5 \cdot 10^{13}$	1	$2 \cdot 10^8$	$8 \cdot 10^{27}$	fission reactions
<sup>238</sup> U <sup>72+</sup>	10	72+	$2 \cdot 10^8$	H <sub>2</sub> , D <sub>2</sub>	$3 \cdot 10^{13}$	1.5	$2 \cdot 10^8$	$5 \cdot 10^{27}$	
<sup>232</sup> Th <sup>71+</sup>	10	71+	$2 \cdot 10^8$	He	$3 \cdot 10^{12}$	1	$2 \cdot 10^8$	$4 \cdot 10^{26}$	
<sup>238</sup> U <sup>72+</sup>	10	72+	$2 \cdot 10^8$	He	$2 \cdot 10^{12}$	1	$2 \cdot 10^8$	$3 \cdot 10^{26}$	

injection charge state

green: bare ion

black: equilibrium charge state after HIE-ISOLDE stripper

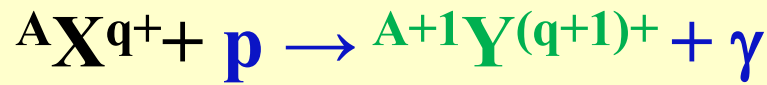
$N_s$ -space charge limit of an electron cooled stored ion beam

required injection rate (ions/s) for  $N_0 = N_s$   
 $n_t$ -target thickness calculated for a beam life time:  $\tau = 1-1.5\text{s}$

remark: H<sub>2</sub>, D<sub>2</sub>, He target: limitation of  $n_t \approx 10^{14}$  atoms/cm<sup>2</sup>

all required injection rate  $N_i/T$  for the space charge limit  $N_s$  maybe can not be reached

# Proton capture reaction for the astrophysical p-process



${}^A\text{X}^{q+}$  - stored main ion

$q$ - charge state main beam

$A$ - mass number

${}^{A+1}\text{Y}^{(q+1)+}$  -daughter nuclide

$p$ - proton from hydrogen target

## 1. Nuclear reactions

momentum conservation

$$A m_0 v_p = (A+1) m_0 v$$

$$\Rightarrow v = \frac{A v_p}{(A+1)}$$

rigidity daughter ion  ${}^{A+1}\text{Y}^{(q+1)+} \Rightarrow B\rho = \frac{p}{Q} = \frac{A}{(q+1)e_0} m_0 v_p$

2. Ionization projectile:  ${}^A\text{X}^{q+} \rightarrow {}^A\text{X}^{(q+1)+} + e$

rigidity stripped ion  ${}^A\text{X}^{(q+1)+} \quad B\rho = \frac{p}{Q} = \frac{A}{(q+1)e_0} m_0 v_p$

daughter  
nuclide

stripped  
ion

same rigidity !!!!!

$\Rightarrow$  rigidities of  ${}^A\text{X}^{(q+1)+}$  and  ${}^{A+1}\text{Y}^{(q+1)+}$  are equal

$\Rightarrow {}^A\text{X}^{(q+1)+}$  and  ${}^{A+1}\text{Y}^{(q+1)+}$  can not separated with magnetic fields !

$\Rightarrow {}^A\text{X}^{(q+1)+}$  and  ${}^{A+1}\text{Y}^{(q+1)+}$  ions are at same detector position

$\Rightarrow$  a) experiment has to carry out with bare  ${}^A\text{X}^{q+}$  ions

$\Rightarrow$  b) bare ions for heavy systems  ${}^A\text{X}^{q+}$  cannot produce by stripping at 10 MeV/u !

looking for solutions to separate  ${}^A\text{X}^{(q+1)+}$  from  ${}^{A+1}\text{Y}^{(q+1)+}$ :

accumulate  ${}^A\text{X}^{(q+1)+}$  in the storage ring and separation with electron cooling

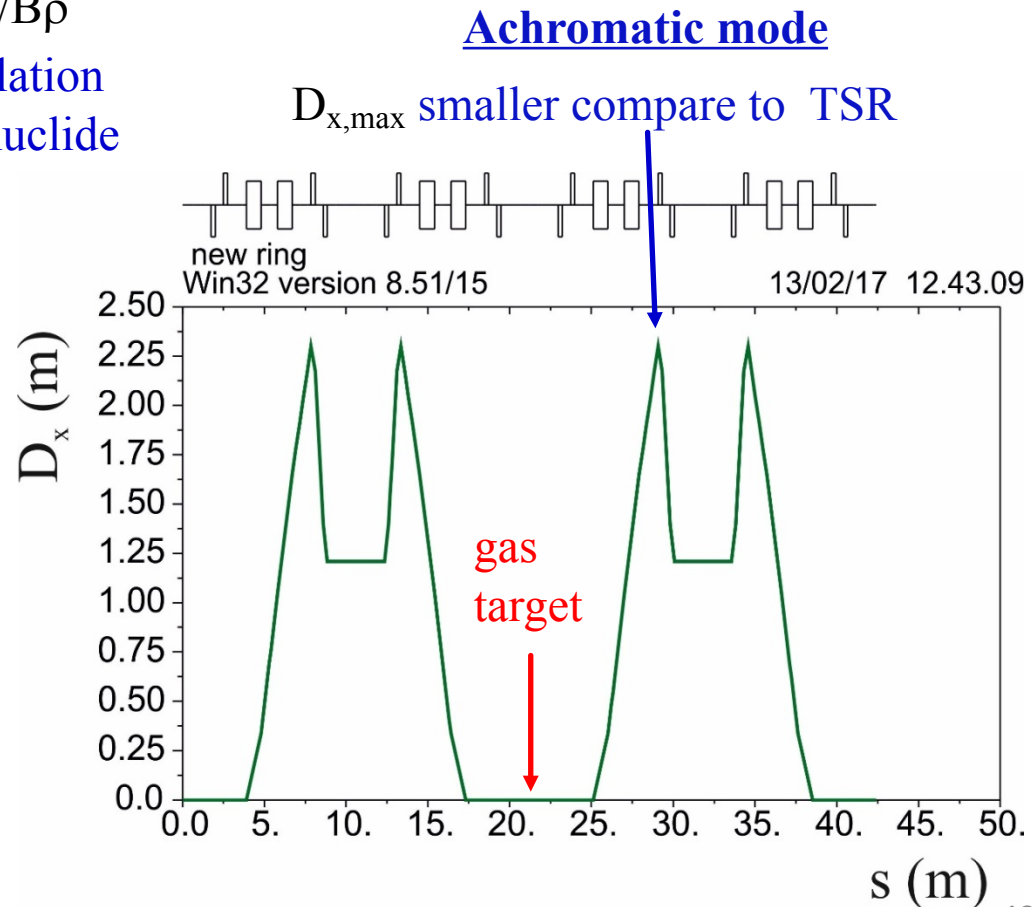
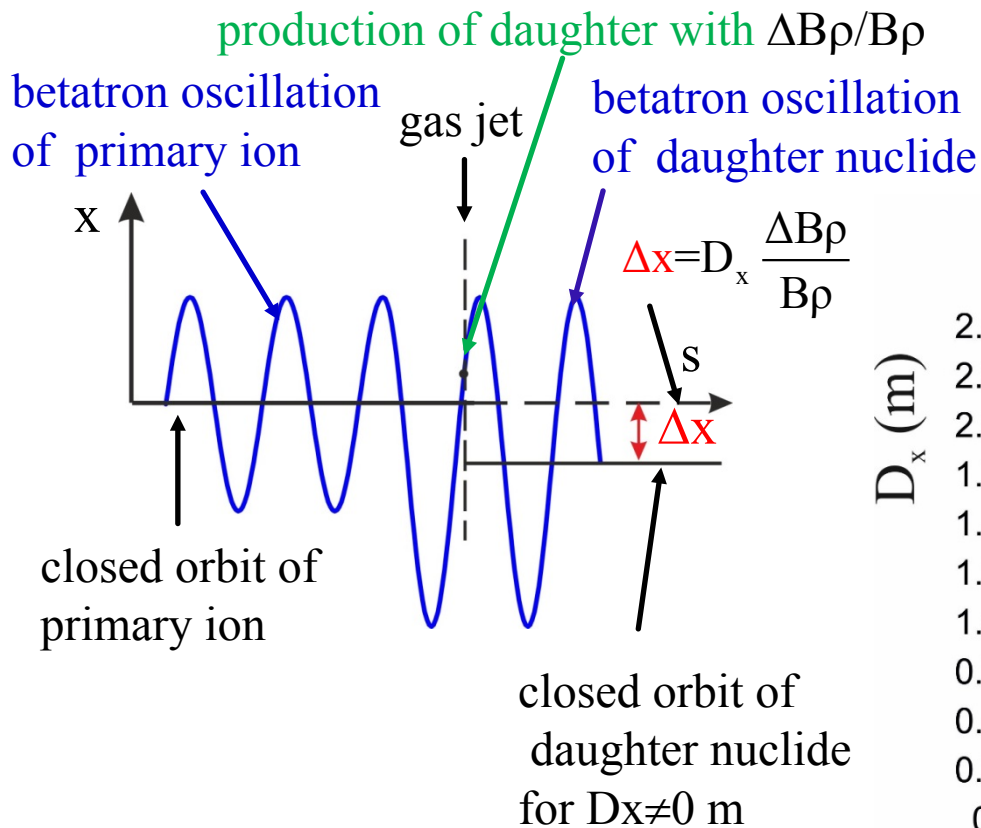


# Proton pick up reactions by storing daughter nuclei

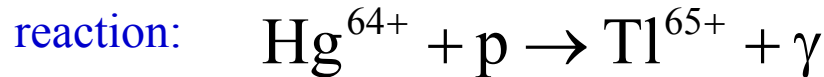
example:  $\text{Hg}^{64+} + p \rightarrow \text{Tl}^{65+} + \gamma$  16 electrons are left !

daughter nuclei  $\text{Tl}^{65+}$  should be kept and accumulated in the storage ring

$\Rightarrow$  storage ring has to operate in an achromatic mode with  $D_x=0$  m in the gas target to avoid excitation of betatron oscillations of the daughter nuclei:



# Shift of the ion orbits by electron cooling



direct after injection:  $v(\text{Tl}^{65+}) < v_e$      $v(\text{Hg}^{65+}) = v_e$

with electron cooling:  $v(\text{Tl}^{65+}) \rightarrow v_e$      $v(\text{Hg}^{65+}) = v_e$

change of the rigidity of  $\text{Tl}^{65+}$ :  $\frac{\Delta B\rho}{B\rho} = -\frac{1}{q+1} \rightarrow -\frac{A-q}{A(1+q)}$

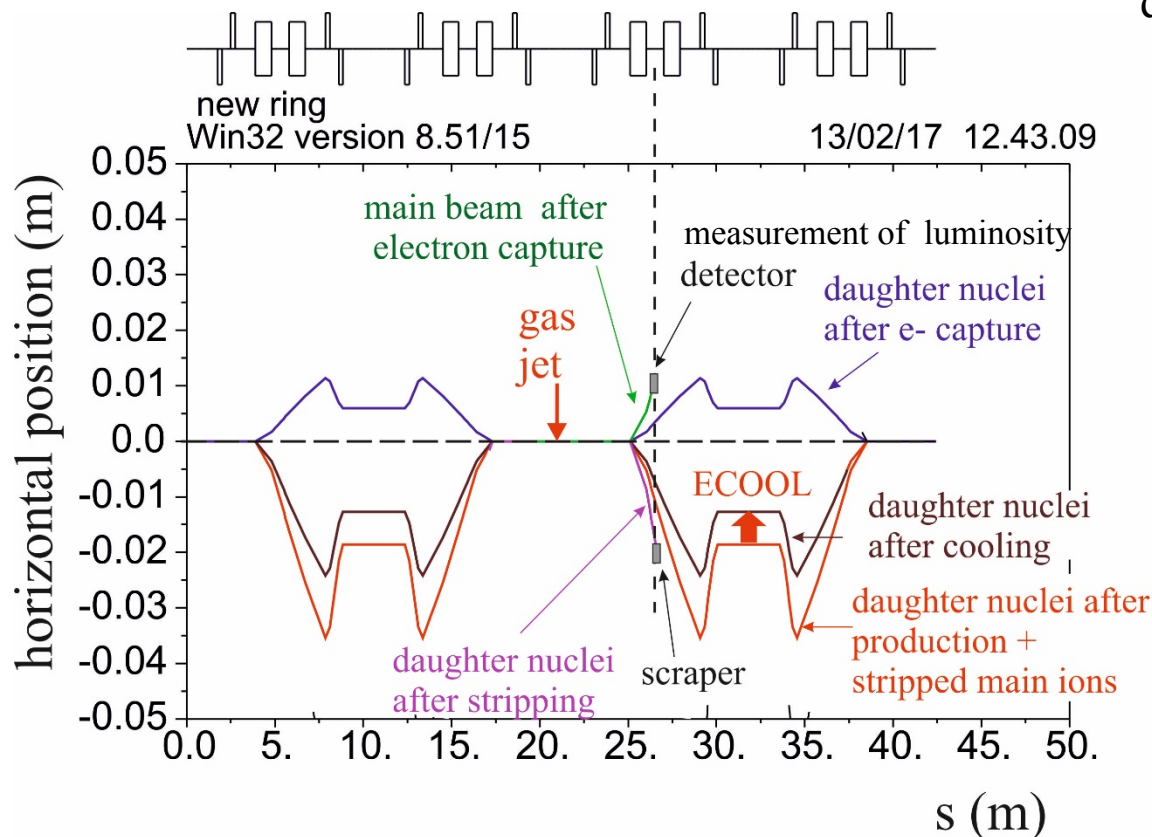
$v_e$  - electron velocity

$v(\text{Hg}^{65+})$  - velocity of  $\text{Hg}^{65+}$

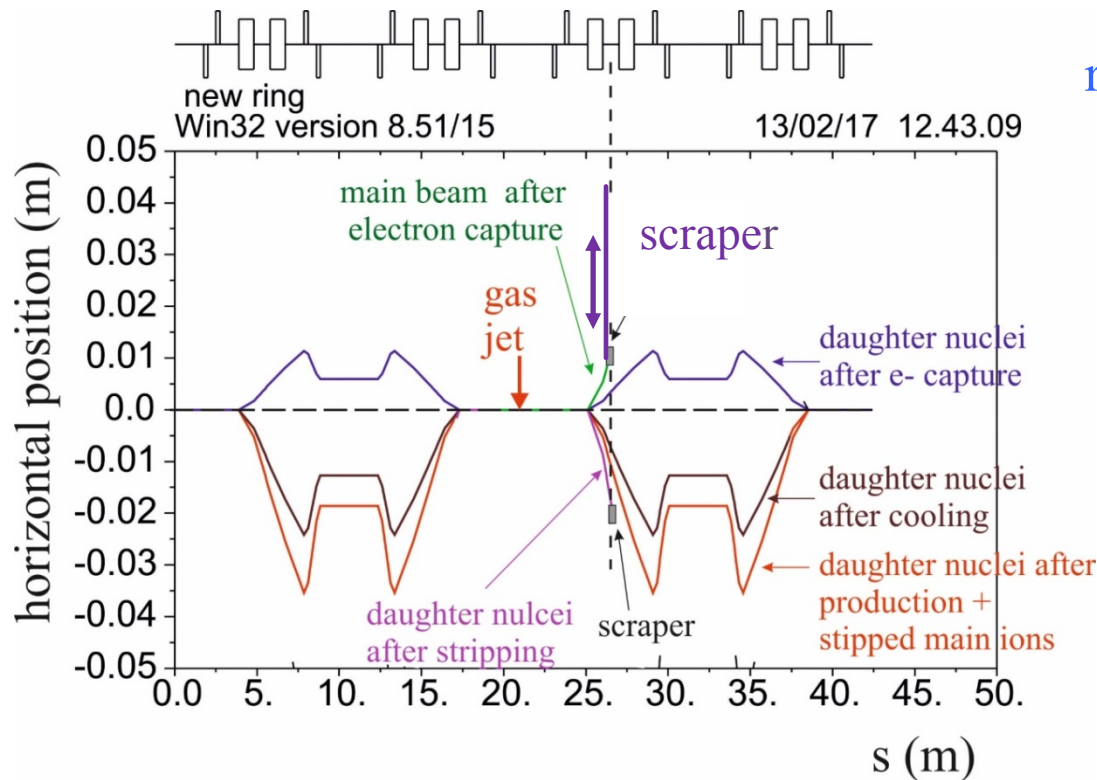
$v(\text{Tl}^{65+})$  - velocity of  $\text{Tl}^{65+}$

A- mass of main beam

q- charge of main beam



# Scraping of electron captured daughter nuclides



rate equation to determine  $N_d$

$$\frac{dN_d(t)}{dt} = L \cdot \sigma - \frac{N_d(t)}{\tau_d}$$

in the equilibrium  
cross section

$$\sigma = \frac{N_{d,0}}{L\tau_d}$$

number of accumulated daughter nuclei

total life time of daughter nuclei

example

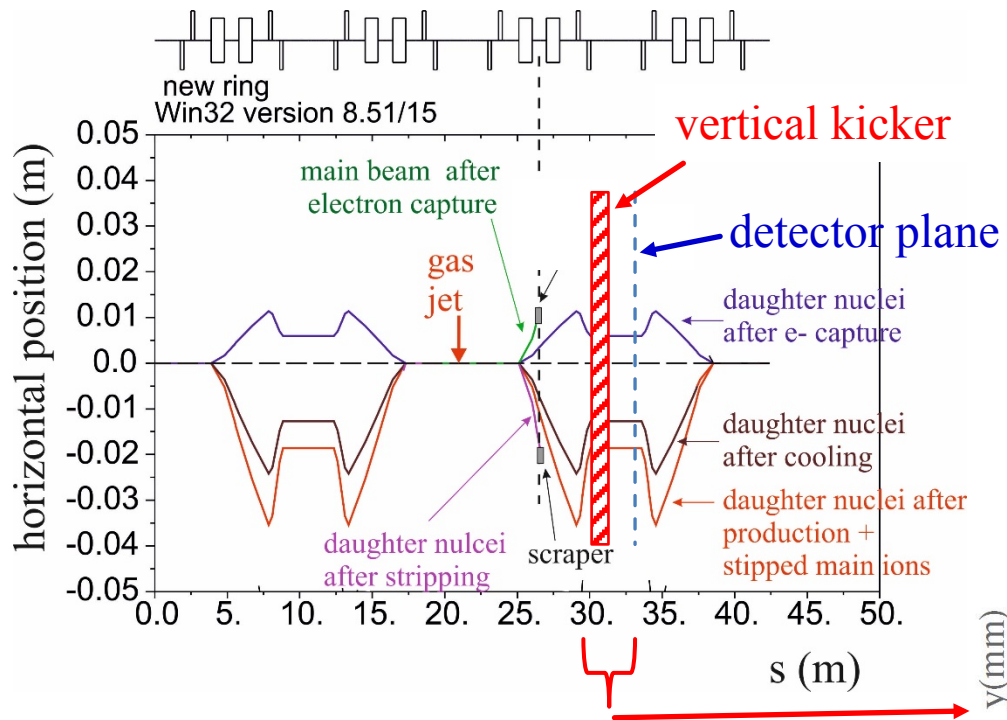
life time should not depend of resonances

beam:  $^{196}\text{Hg}^{64+}$ ,  $E/A=10$  MeV/u,  $N_0 \approx N_s/6 = 3 \cdot 10^7$ ,  $n_t = 2 \cdot 10^{13}$  atoms/cm<sup>2</sup>

$$\Rightarrow L = 2.2 \cdot 10^{27} \text{ 1/cm}^2\text{s}, \tau_d \approx 3.5 \text{ s}$$

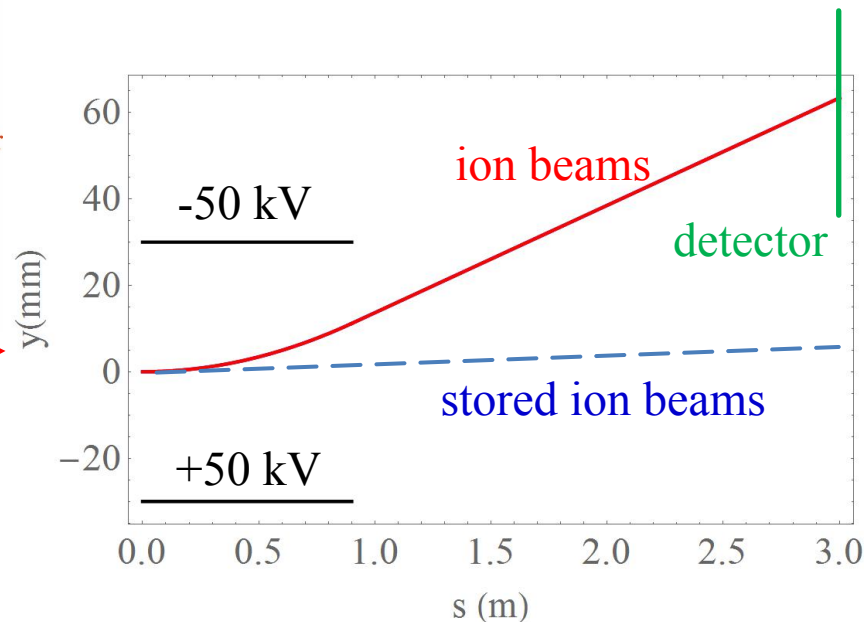
for:  $\sigma = 0.01$  barn :  $N_{d,0} = \sigma L \tau_d \Rightarrow N_{d,0} \approx 80$

# Direct Detection of the produced daughter nuclei



A vertical kicker is used to kick the stored ion beam towards the detectors

## second experimental straight section II



After reaching the equilibrium daughter

nuclei number given by:  $N_{d,0} = L\tau_d\sigma$

The kicker voltage is switched on very quickly with Behlke switches to deflect the stored ion beam to the detectors located above to the stored ion beam, where the switching time  $T \ll T_0$  ( $T_0$  -revolution time)



# Multi Charge operation of the TSR

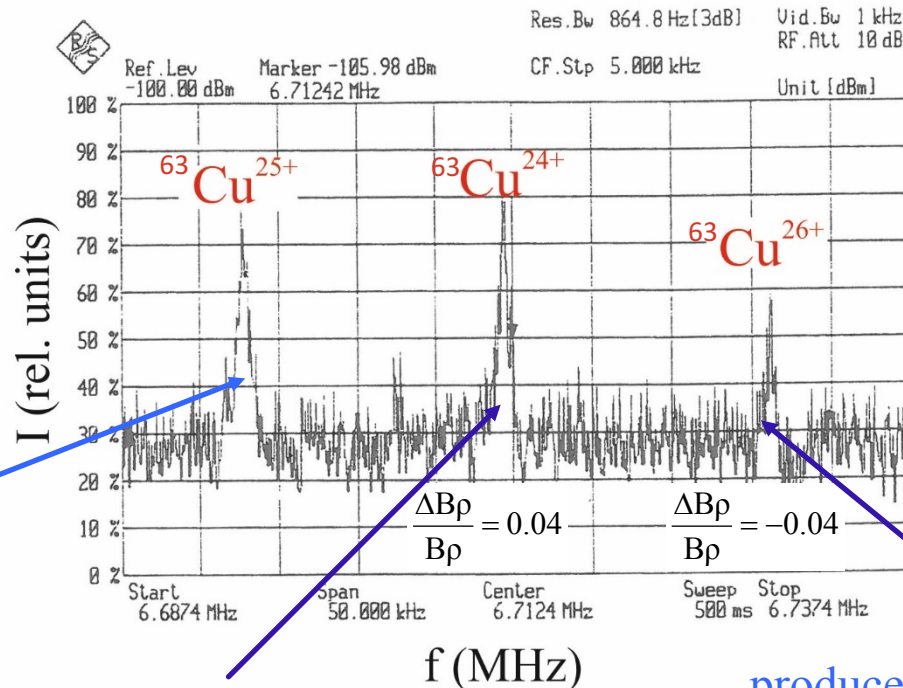
- Storing also daughter nuclei after production needs a relative large momentum acceptance of the storage ring
- At the TSR it was shown that several beams with different rigidities can be stored at the same time

## Confirmation of the multi charge operation at the TSR

Schottky noise measured 12 s after injection

dispersion in the cooler:

$$D_x = 0.3 \text{ m}$$



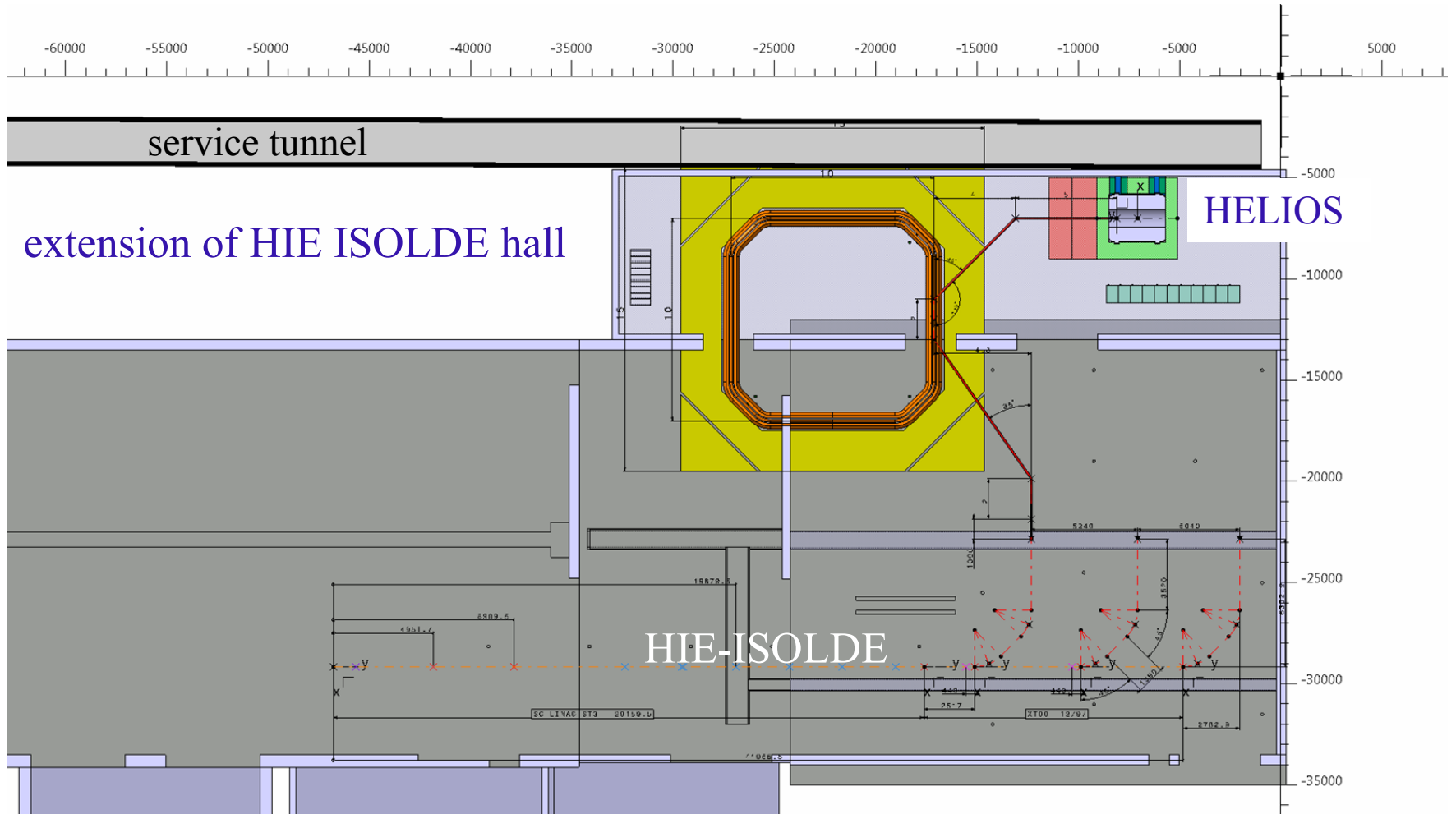
injected ion beam  
 $E = 266 \text{ MeV}$

produced by electron capture  
visible already 2-3 s after injection

produced by stripping  
detectable 8 s after injection

# Possible location of the new storage at HIE-ISOLDE

transparency from Erwin Siesling and Stephane Maridor, CERN



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